



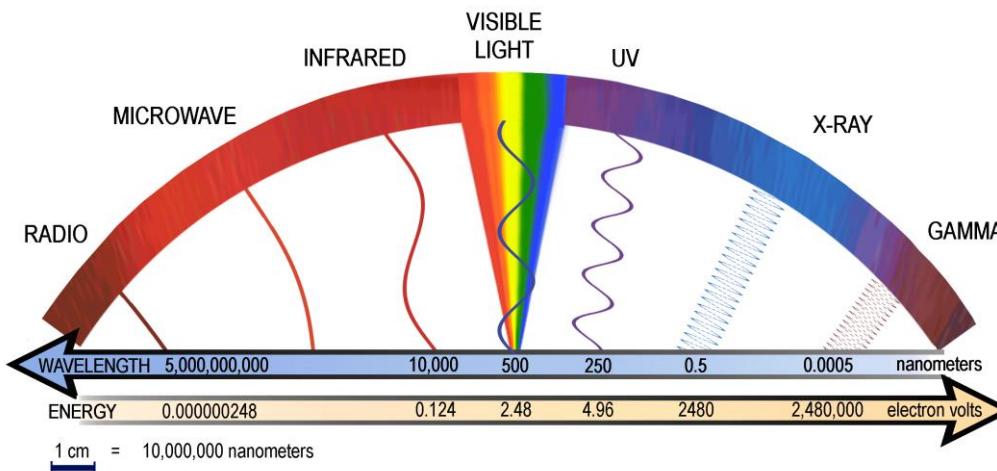
University
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Global energetics, particle acceleration and turbulence in standard solar flare model

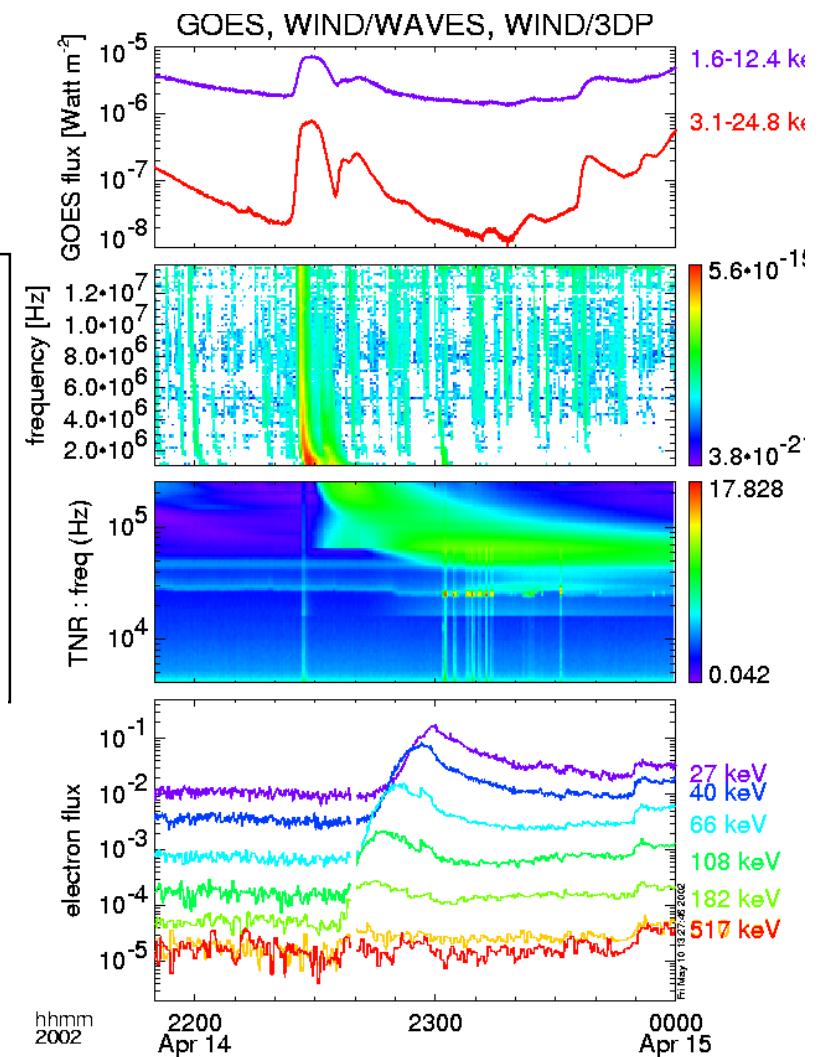
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J. E. Perez, L. K. Harra, A. A. Kuznetsov, A. G.
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Solar flares are rapid localised brightening in the lower atmosphere.

More prominent in X-rays, UV/EUV and radio.... but can be seen from radio to 100 MeV



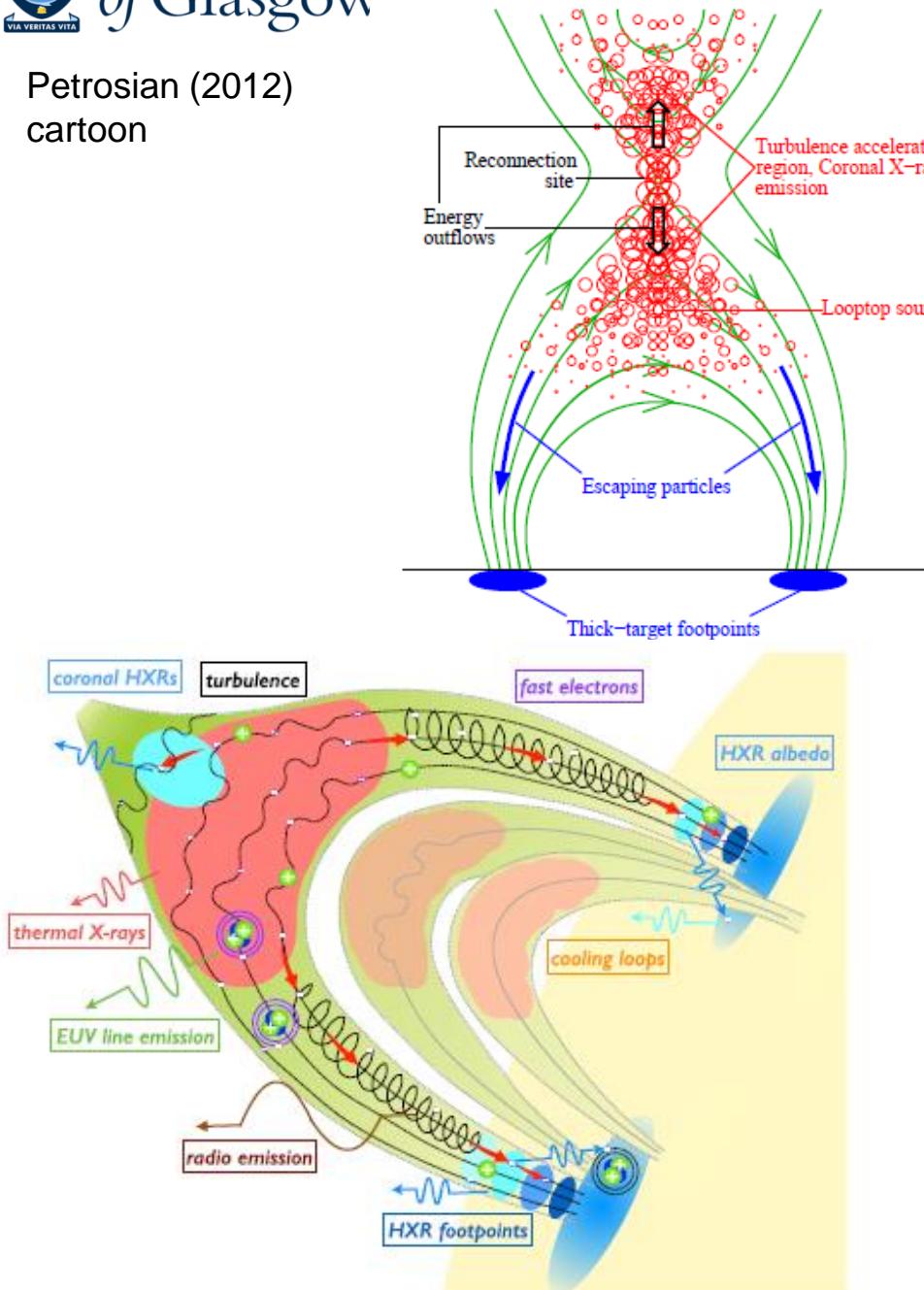
X-rays



Particles 1AU

Figure from Krucker et al, 2007

Petrosian (2012)
cartoon



Standard model energetics

Magnetic Energy



Turbulence/Fluctuating E field



Acceleration/Heating



Electrons/Ions

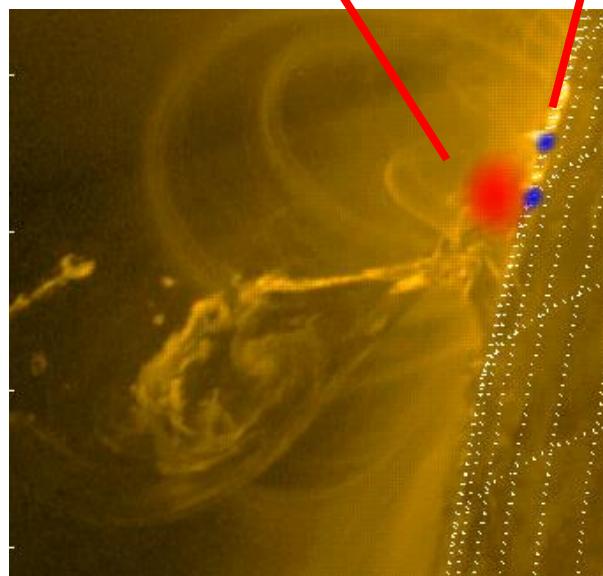
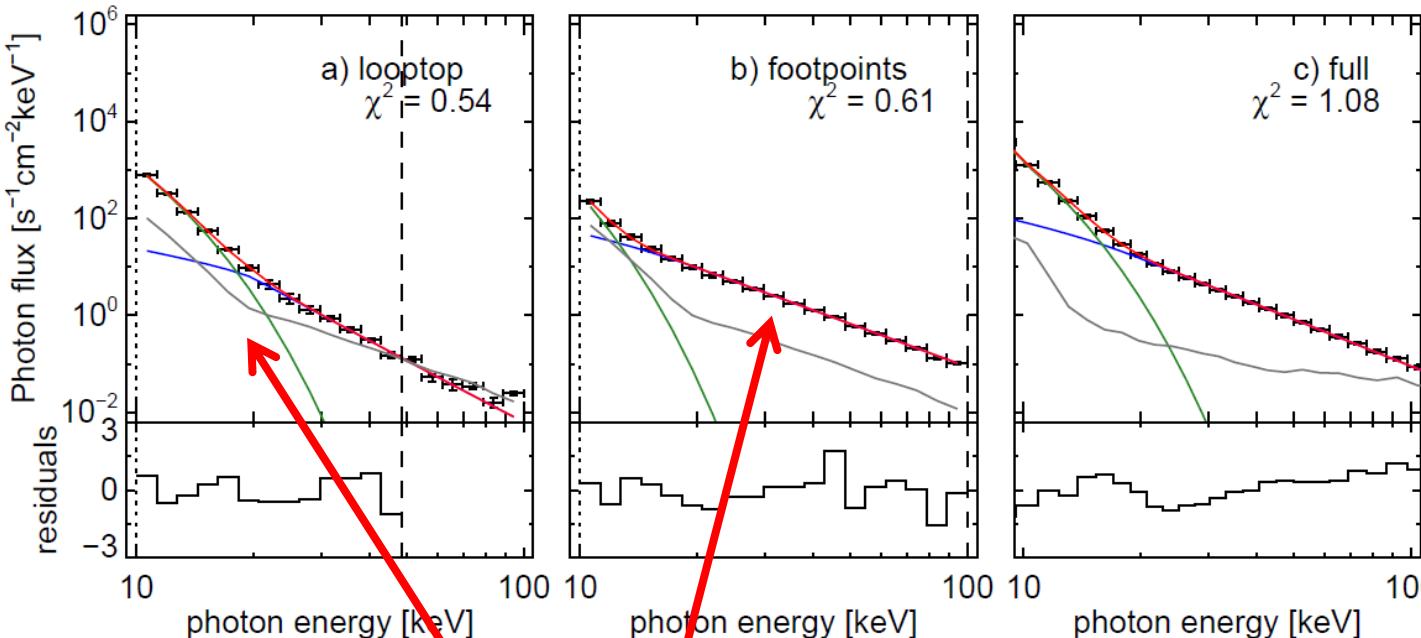


Energy Deposition/Evaporation



Radiation

Typical solar flare: X-ray prospective



Loop-top: Soft X-ray plus non-thermal component

Footpoints: Hard X-ray non-thermal power-law

Simoes & Kontar, A&A, 2013

Using imaging spectroscopy, we can infer spectra and numbers of energetic electrons both in coronal and foot-points sources.

Above 30 keV, we have normally a few times electrons more in the LT than in FP source.
Possible trapping by waves or mirror?

Stochastic particle acceleration

Vast literature exists e.g. Miller et al, 1997; Petrosian 2012; Bian et al, 2012 Cargill et al, 2012 as reviews

⇒ Generally efficient electron and proton acceleration, He₃ enhancement, variety and variability of particle spectra

Particle and energy transport

Pitch angle scattering of particles, reduced thermal conductivity, etc

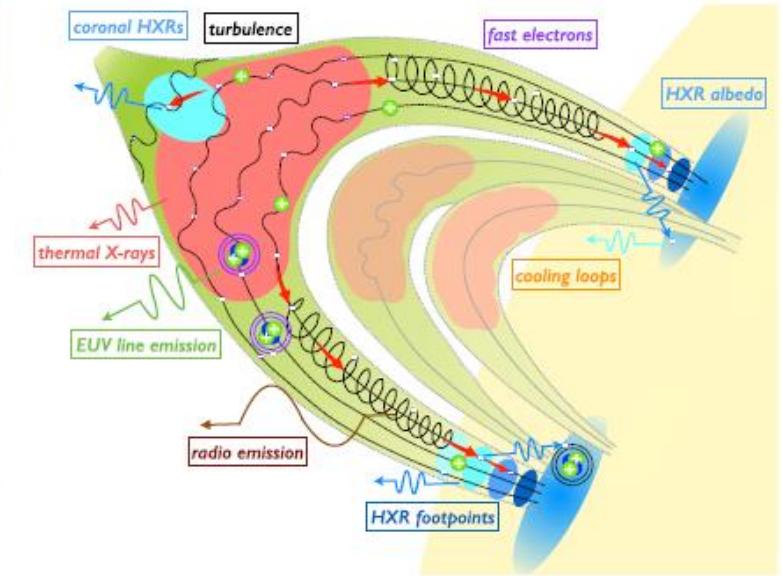
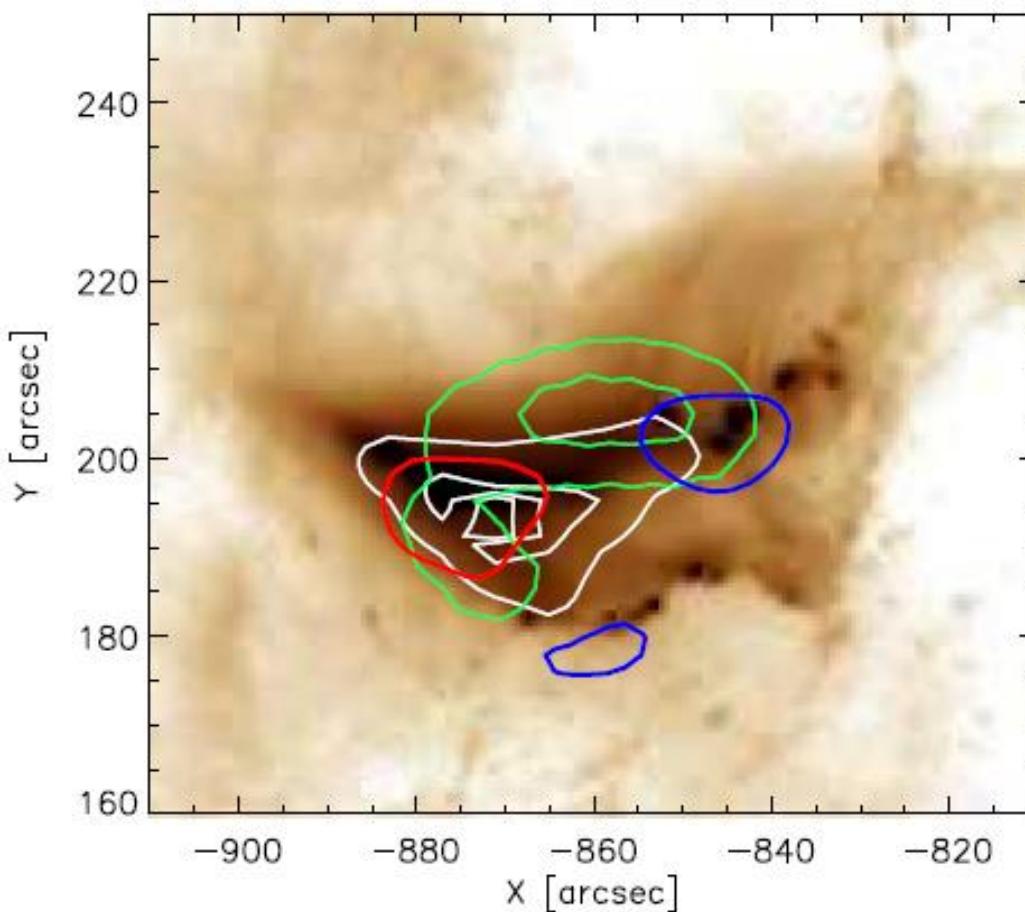
⇒ Artificial injection of electrons often involved to explain strong radio sources at the loop-tops
⇒ (e.g. Melnikov et al 2001, Lee et al, 2002,
⇒ Simoes & Kontar 2013, Kontar et al, 2014)

Reconnection models

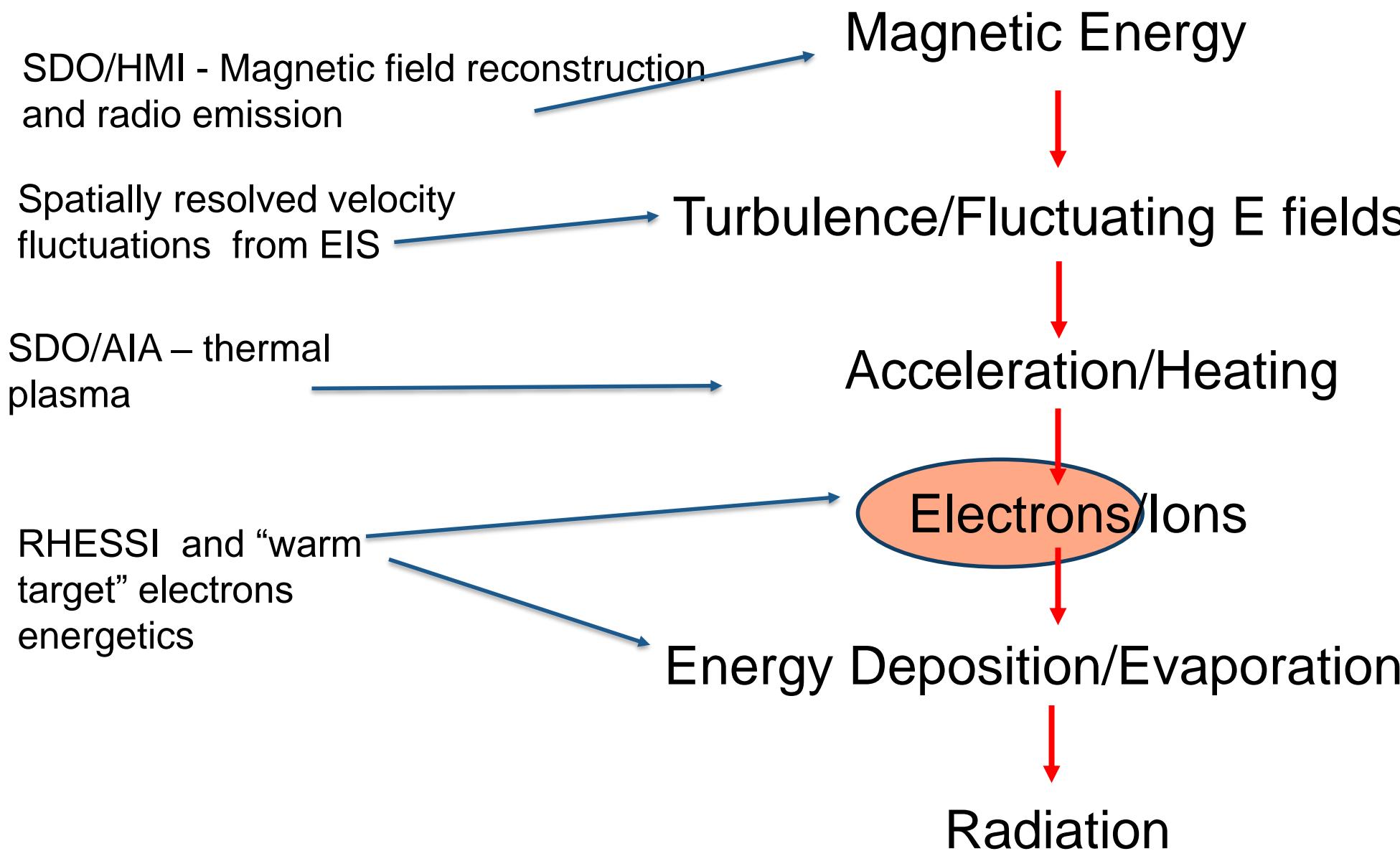
Anomalous resistivity is often required to make fast reconnection and strong parallel electric fields

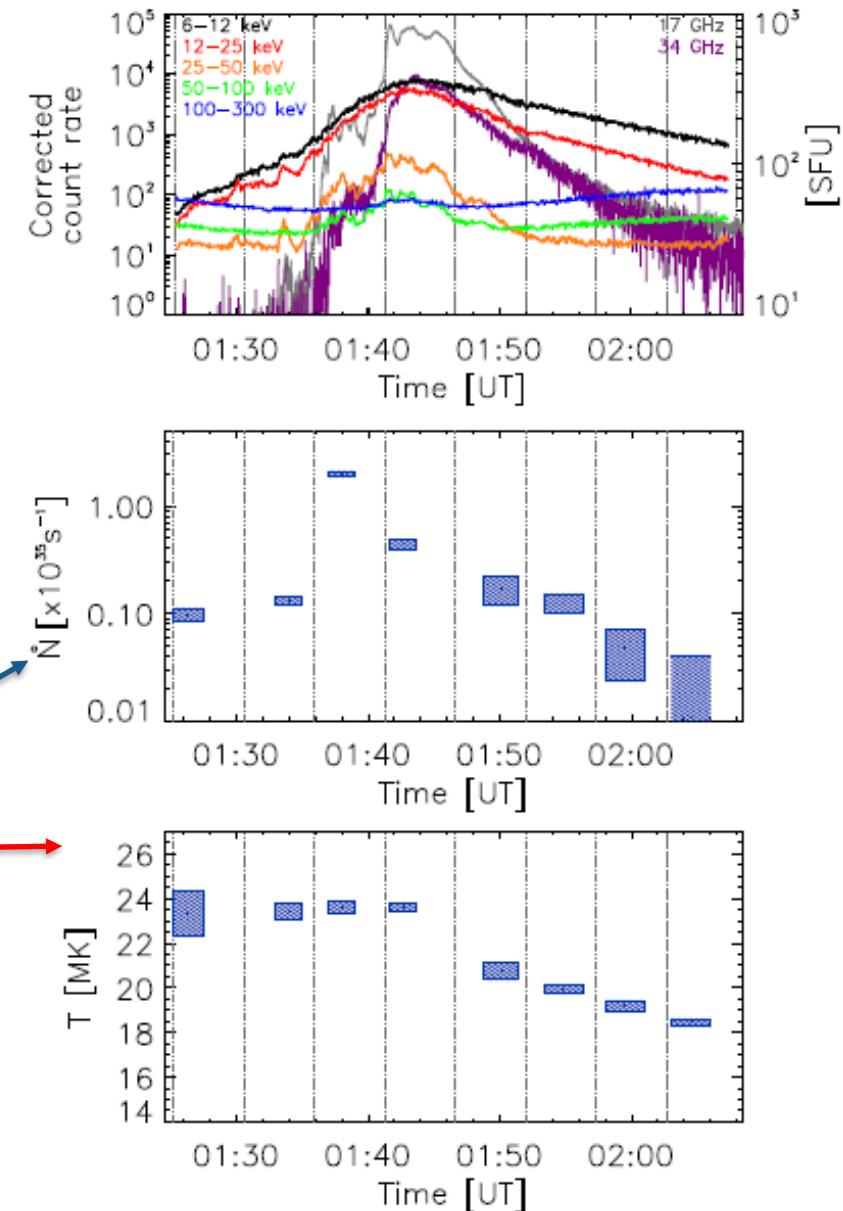
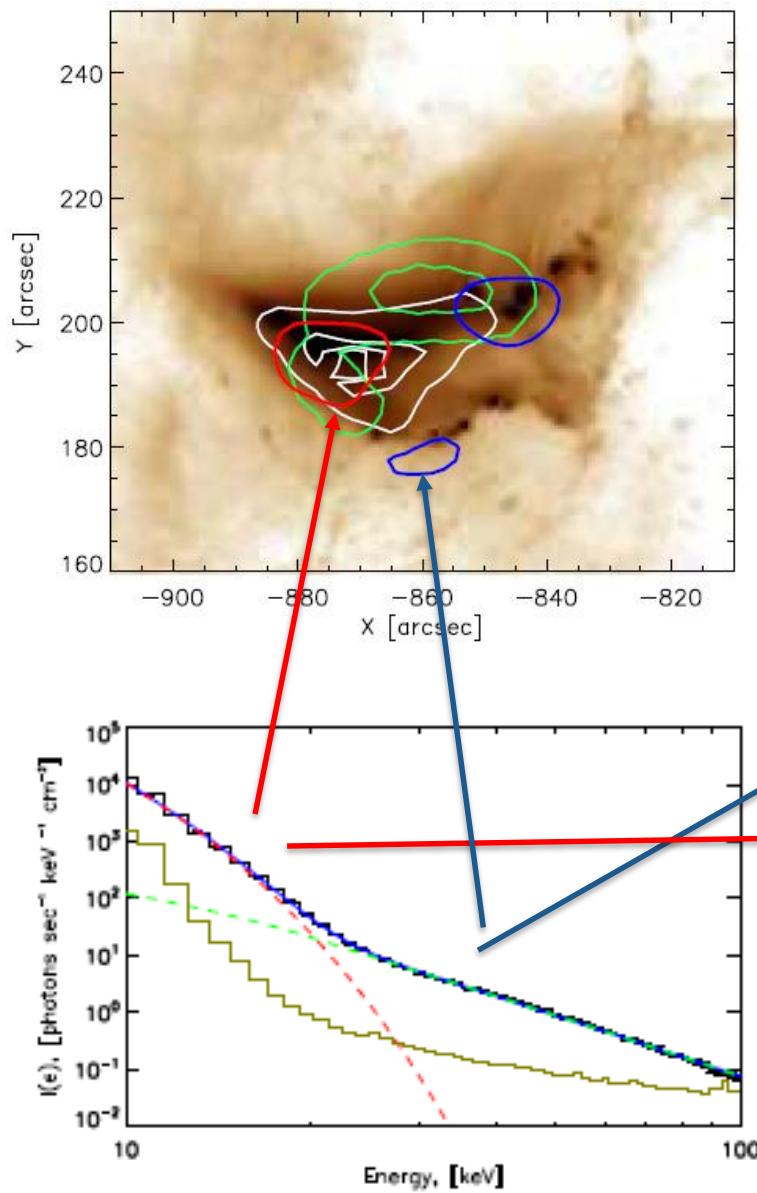
see e.g. Lazarian et al, 1999, Raymond et al, 2012, Gordovskyy et al, 2016

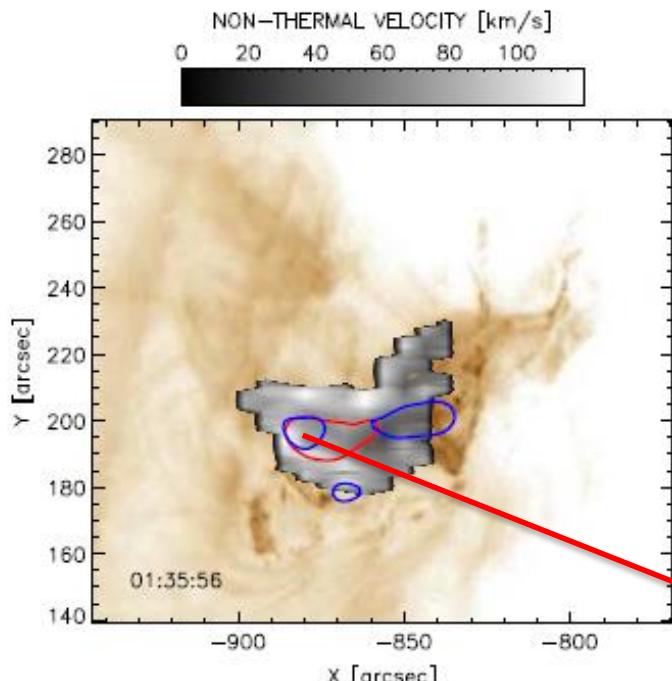
Plasma turbulence is characterised by chaotic and stochastic property changes, velocity, density, magnetic field in space and time.



SDO/Atmospheric Imaging Assembly (AIA) 193 Å image (background); RHESSI x-ray contours at 50% of peak value for 6–15 (red) and 25–50 keV (blue) energy ranges, EIS Fe XXIV (255 Å) intensity map (white contours at 30% and 75% of peak value), and Nobeyama 34 GHz radio emission (green contours at 30% and 75% of peak value).

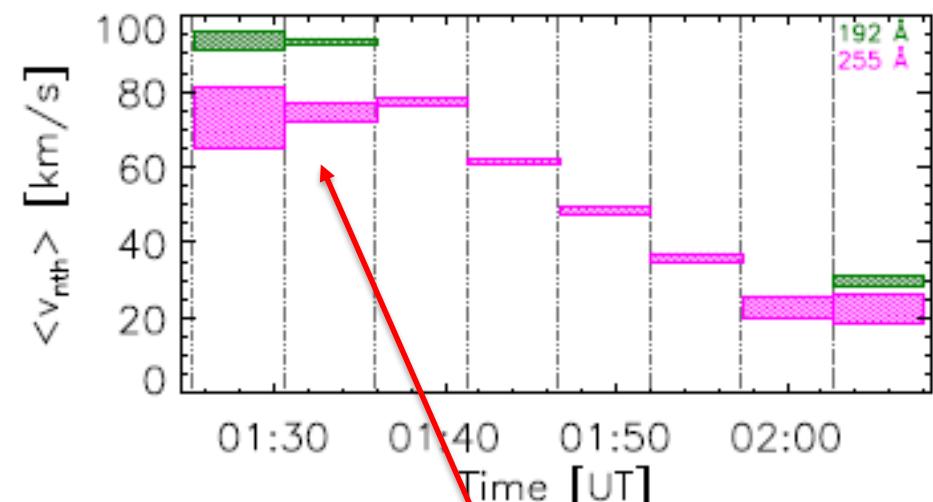






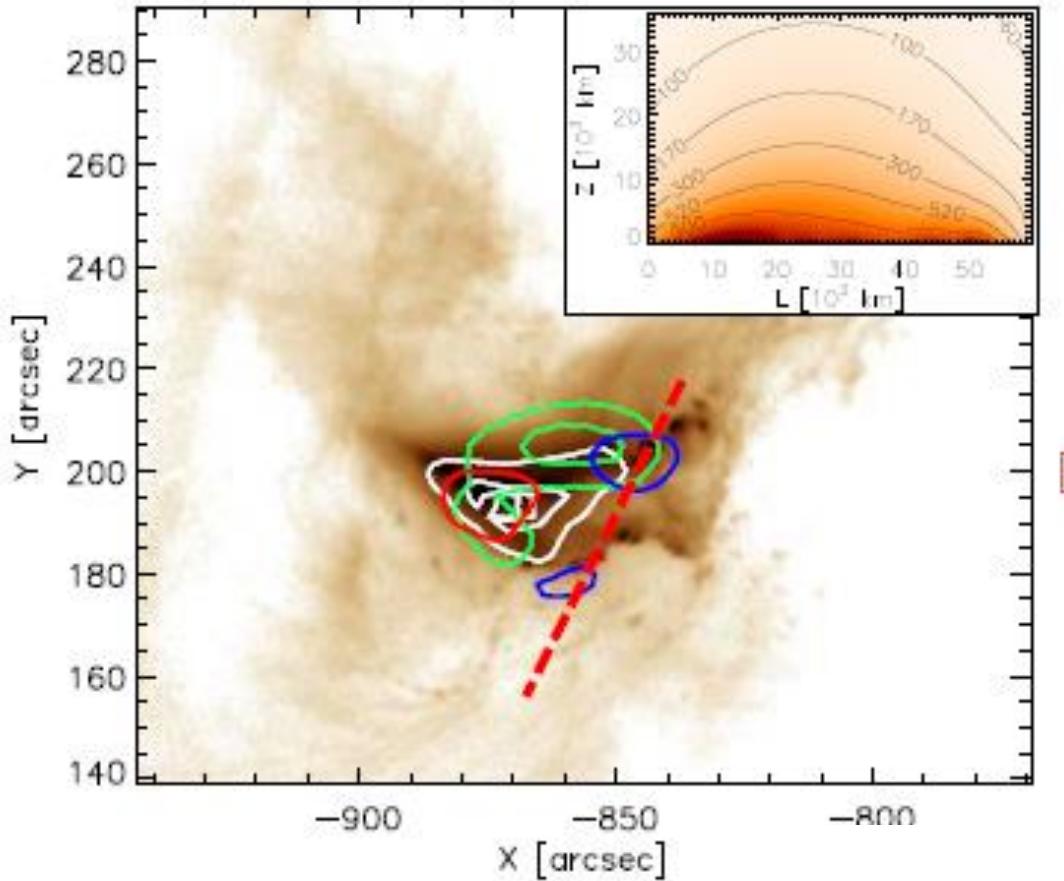
See Harra et al 2013 for details

EIS Line width (Fe XXIV) => plasma velocity variance

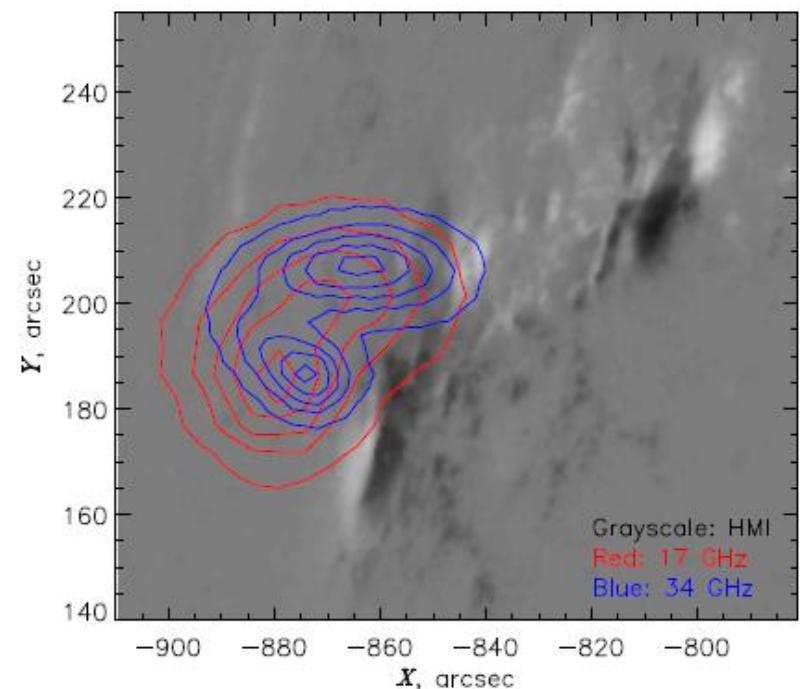
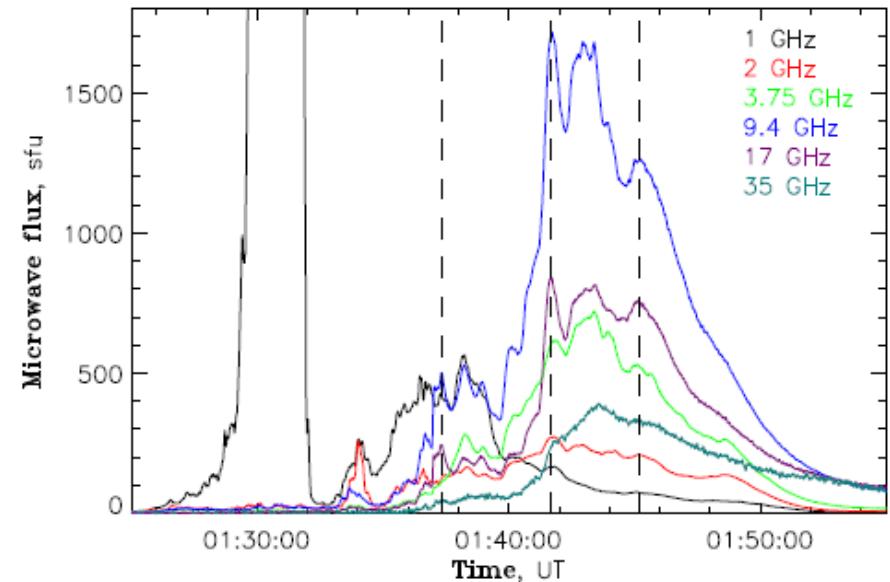


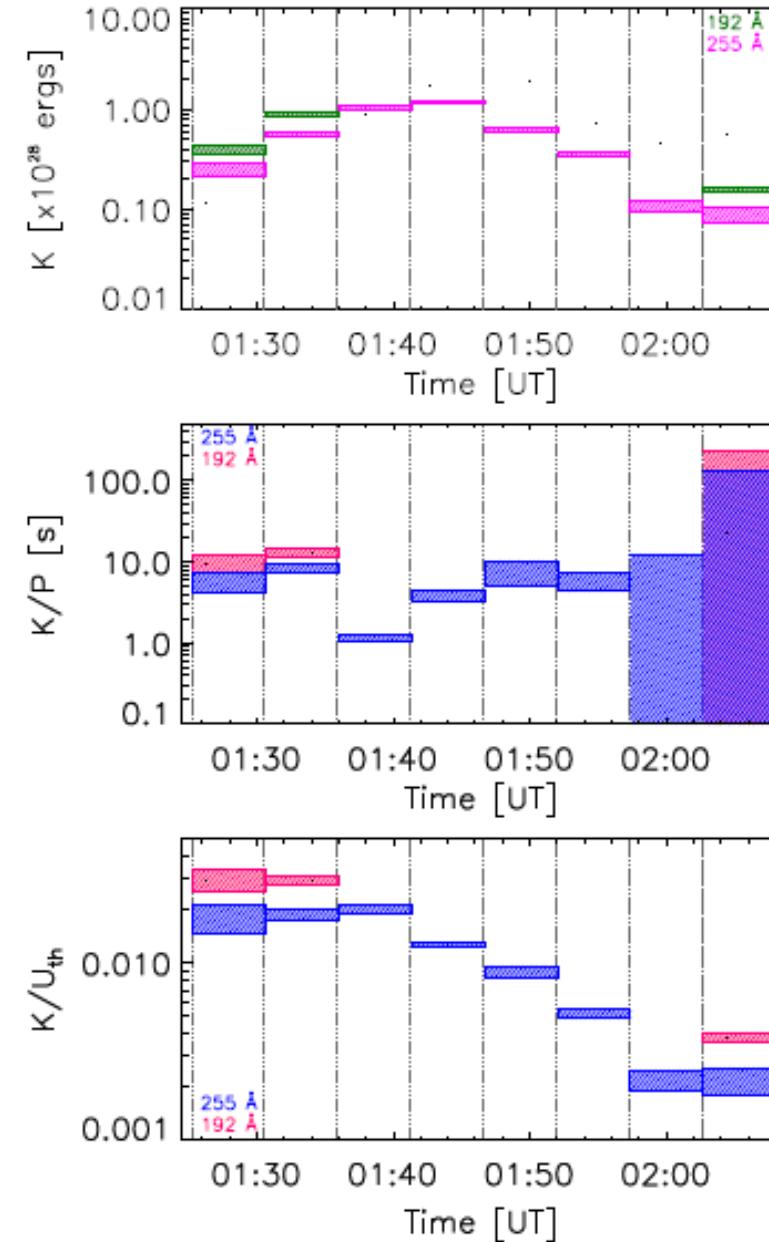
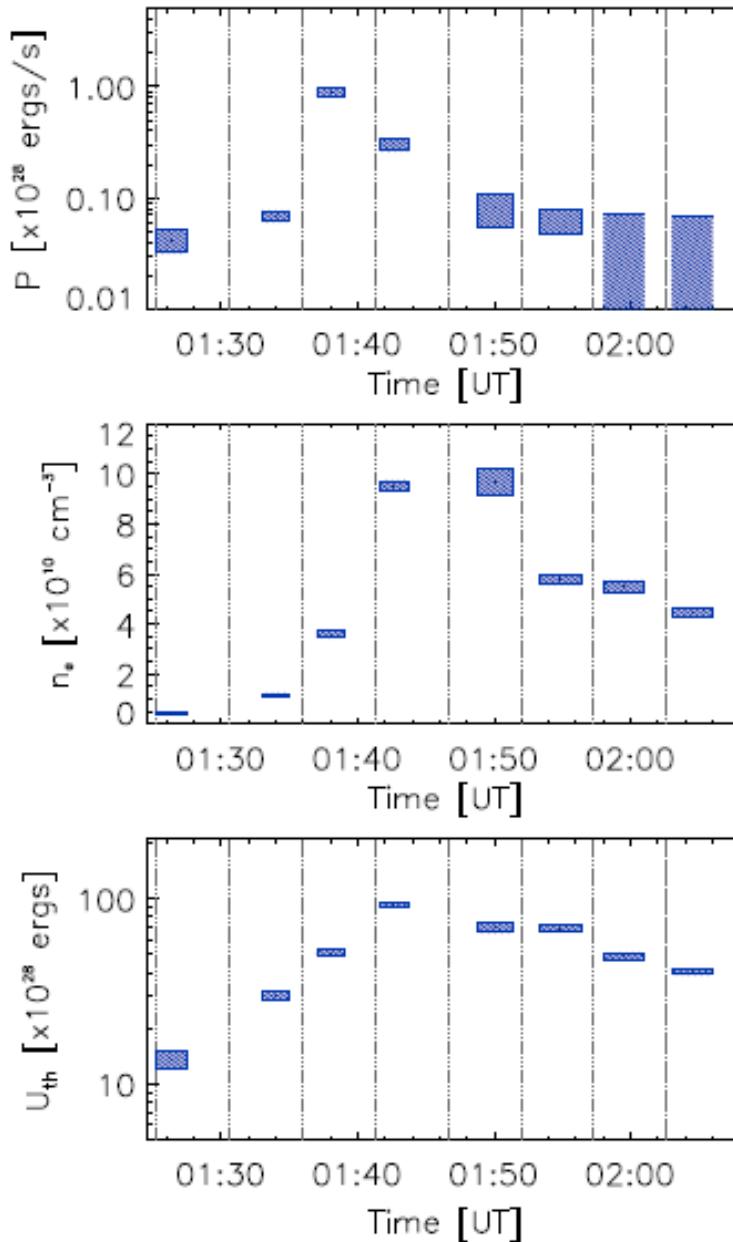
$$\text{FWHM}^2 = (\text{instr}_{\text{fwhm}})^2 + 4 \ln(2)(\lambda/c)^2 (v_t^2 + (v_{\text{nt}})^2)$$

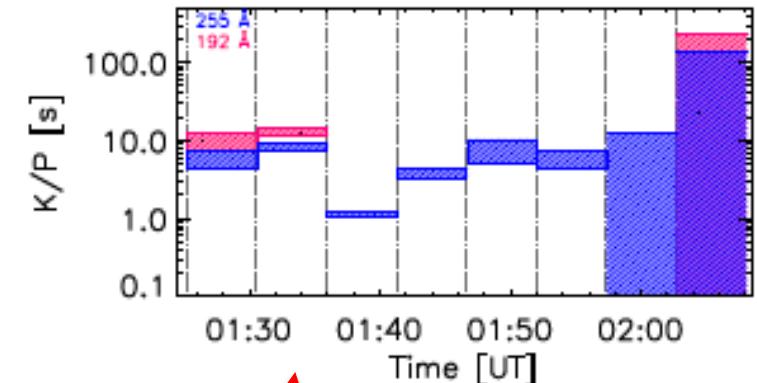
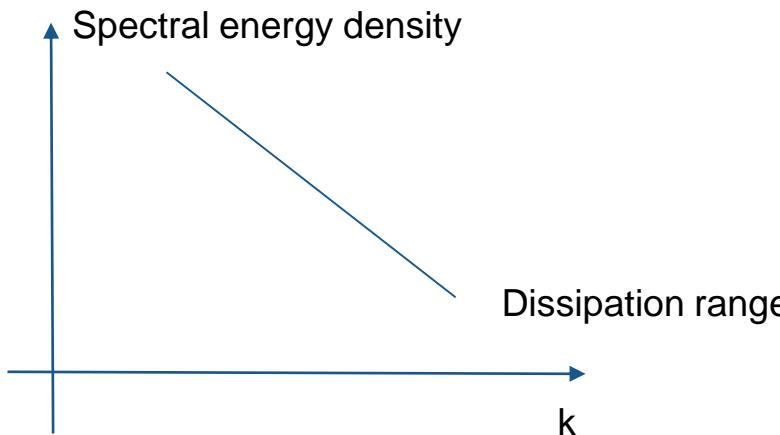
e.g. non-thermal line broadening observed in flares, normally spatially unresolved, but now with Hinode/EIS – spatially resolved



$$U_m = \frac{B^2 L^3}{8\pi}$$



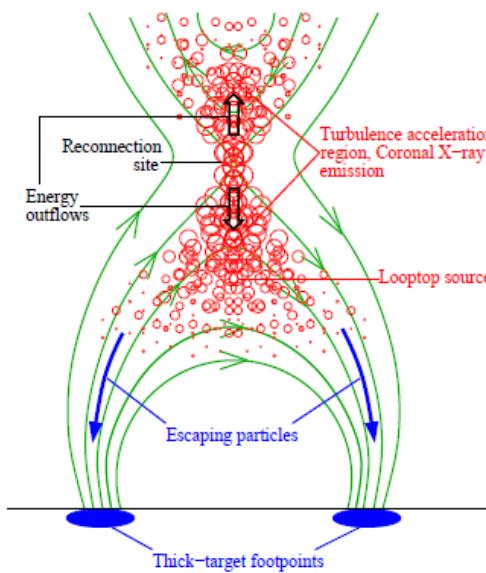




Turbulence dissipation timescale $\sim L_{\perp}/v_{nth}$.

P. Goldreich and S. Sridhar, 1995

The energy density associated with a turbulence-perturbed magnetic field δB is $U_B \simeq (\delta B)^2/8\pi$. Equating this to the turbulent energy content (Alfvénic MHD turbulence) $K = (1/2)n m \langle v_{nth}^2 \rangle$, we obtain $\langle v_{nth}^2 \rangle \simeq (\delta B)^2/4\pi n m$. Since the Alfvén speed $V_A = \sqrt{B^2/4\pi n m}$, it follows that $\langle v_{nth} \rangle/V_A \simeq \delta B/B \simeq L_{\perp}/L_{\parallel}$, where L_{\parallel} is the longitudinal extent of the turbulence region. Thus the dissipation timescale $L_{\perp}/\langle v_{nth} \rangle$ is approximately the same as the Alfvén crossing time L_{\parallel}/V_A , a quantity that is readily ascertainable from observations. Using the values found in the observations, the dissipation of turbulent energy occurs on a time scale $\sim 1 - 10$ seconds.



Magnetic Energy

$\sim 10^{31}$ ergs

Turbulence

$\sim 10^{28}$ ergs

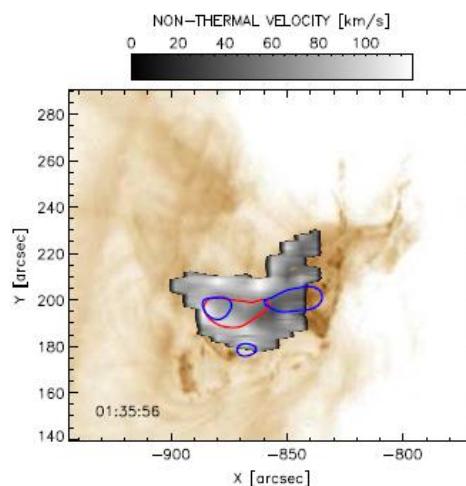
Electron Acceleration

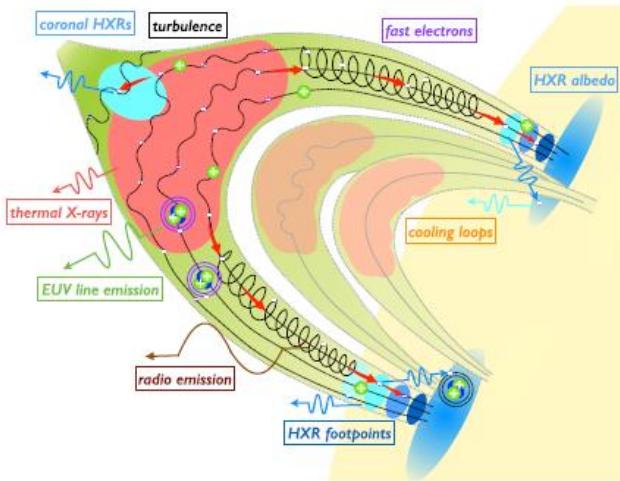
$\sim 10^{28}$ erg/sec

Thermal energy

$\sim 10^{30}$ ergs

Radiation





The new multiwavelength observations provide persuasive evidence that plasma turbulence plays a key role in the energy transfer in solar flares, thus directly supporting the stochastic acceleration model.

The turbulent kinetic energy (instantaneous) comprises only a small part of the total flare energy budget ($\sim 0.2\text{-}1\%$). Nevertheless, provided that its energization and dissipation processes are rapid enough (with the timescales of $\sim 1\text{-}10$ s),

The energy dissipation rate and the power in accelerated nonthermal particles observed in the flare are consistent with the dissipation of anisotropic Alfvén MHD turbulence.

The observations could be used to test various acceleration models.