

Accurate theory of the hydrogen molecule

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Motivation

- benchmark calculation of relativistic and QED effects in molecular systems
- improved comparison with experimental values for H_2 levels (currently at 10^{-8} precision level)
- m_e/m_p from the vibrational spectrum, other constants ...
- to discover new effects, like that in muonic versus electronic-hydrogen Lamb shift (PSI 2010), any hints ?

Collaborators

- J. Komasa (Poznań, Poland)
- B. Jeziorski and his group (Warsaw, Poland)

Determination of the dissociation energy of H_2 – a historical perspective

- 1927 2.88 eV Heitler and London (lower bound)
- 1929 4.46(4) eV Richardson and Davidson present value 4.478 eV
- 1935 4.454(13) eV James and Coolidge
- 1960 36113.6(3) Herzberg and Monfils
- 1968 36117.4 Kolos and Wolniewicz (lower bound)
- 1970 36118.6(5) Stwalley
- 1993 36118.049 Kolos and Rychlewski
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- 2004 36118.062(10) Zhang *et al.* (Eyler group)
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Challenges

- beyond Born-Oppenheimer: combined finite nuclear mass and (relativistic) QED corrections to rovibrational energies, shielding constant,...
- higher order QED effects $\sim \alpha^4 Ry$
- extension of the computational technique to an arbitrary light molecule

Atomic or molecular energy

$\alpha \approx 1/137$ – fine structure constant, $\beta = m_e/\mu_n \sim 10^{-3}$ electron-nuclear mass ratio

$$\begin{aligned}
 E(\alpha, \beta) = & E_0 + \alpha^2 E_{\text{REL}} + \alpha^3 E_{\text{QED}} + \alpha^4 E_{\text{HQED}} + \dots \\
 & + \qquad \qquad \qquad + \qquad \qquad \qquad + \qquad \qquad \qquad \vdots \\
 & \beta E_{\text{AD}} \qquad \alpha^2 \beta E_{\text{REL}}^{\text{rec}} \qquad \alpha^3 \beta E_{\text{QED}}^{\text{rec}} \\
 & + \qquad \qquad \qquad \vdots \qquad \qquad \qquad \vdots \\
 & \beta^2 E_{\text{NA}} \\
 & \vdots
 \end{aligned}$$

E_0 – Born-Oppenheimer energy

E_{AD} – adiabatic correction

E_{NA} – nonadiabatic correction

$\alpha^2 E_{\text{REL}}$ – relativistic correction

$\alpha^3 E_{\text{QED}}$ – leading radiative (QED) correction

$\alpha^4 E_{\text{HQED}}$ – higher order QED correction

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Nonrelativistic Schrödinger equation

$$[H - E] |\phi\rangle = 0$$

$$H = H_{\text{el}} + H_{\text{n}}$$

$$H_{\text{el}} = - \sum_a \frac{\nabla_a^2}{2m_e} + V$$

$$H_{\text{n}} = - \frac{\nabla_{\vec{R}}^2}{2\mu_n} - \frac{\nabla_{\text{el}}^2}{2\mu_n} - \left(\frac{1}{M_B} - \frac{1}{M_A} \right) \vec{\nabla}_R \cdot \vec{\nabla}_{\text{el}}$$

$$\vec{R} = \vec{R}_A - \vec{R}_B$$

$$\frac{1}{\mu_n} = \frac{1}{M_B} + \frac{1}{M_A}$$

$$\vec{\nabla}_{\text{el}} \equiv \frac{1}{2} \sum_a \vec{\nabla}_a$$

Diatomic molecules, coordinate system origin at the bond midpoint

Adiabatic approximation

perturbative expansion

$$\phi = \phi_a + \delta\phi_{\text{na}}$$

$$E = E_a + \delta E_{\text{na}}$$

adiabatic assumption

$$\phi_a(\vec{r}, \vec{R}) = \phi_{\text{el}}(\vec{r}) \chi(\vec{R})$$

$$[H_{\text{el}} - \mathcal{E}_{\text{el}}(\vec{R})] |\phi_{\text{el}}\rangle = 0$$

nuclear equation

$$[H_n + \mathcal{E}_{\text{el}}(R) + \mathcal{E}_a(R) - E_a] |\chi\rangle = 0$$

$$\mathcal{E}_a(R) = \langle \phi_{\text{el}} | H_n | \phi_{\text{el}} \rangle_{\text{el}}$$

Nonadiabatic corrections from NAPT

Comparison with direct variational calculations

Molecule	Method	E_{00}
H ₂	Variational ¹	36 118.79756(2)
	NAPT ²	36 118.7978(2)
	Difference	0.0002
D ₂	Variational	36 749.0910
	NAPT ³	36 749.0910(2)
	Difference	0.0000
HD	Variational ⁴	36 406.5105(4)
	NAPT ⁵	36 406.5108(2)
	Difference	0.0003

¹M. Stanke, D. Kędziera, S. Bubin, M. Molski, L. Adamowicz, *JCP* **128**, 114313 (2008);

²K. Pachucki, J. Komasa, *JCP* **130**, 164113 (2009);

³Piszczatowski, Lach, Przybytek, Komasa, Pachucki, Jeziorski, *JCTC* **5**, 3039 (2009)

⁴M. Stanke, S. Bubin, M. Molski, L. Adamowicz, *Phys. Rev. A* **79**, 032507 (2009);

⁵K. Pachucki, J. Komasa, *PCCP* **12**, 9188 (2010)

Relativistic correction

$$E(\alpha) = (E_0 + E_{\text{AD}} + E_{\text{NA}}) + \alpha^2 E_{\text{REL}} + \alpha^3 E_{\text{QED}} + \alpha^4 E_{\text{HQED}}$$

Expectation value of the Breit-Pauli Hamiltonian

$$\begin{aligned} \mathcal{E}_{\text{REL}} &= \sum_a \left[-\frac{1}{8} \langle \phi_{\text{el}} | \mathbf{p}_a^4 | \phi_{\text{el}} \rangle + \sum_I \frac{Z_I \pi}{2} \langle \phi_{\text{el}} | \delta(\vec{r}_{aI}) | \phi_{\text{el}} \rangle \right] \\ &+ \sum_{a>b} \left[\pi \langle \phi_{\text{el}} | \delta(\vec{r}_{ab}) | \phi_{\text{el}} \rangle + \frac{1}{2} \left\langle \phi_{\text{el}} \left| \vec{p}_a \frac{1}{r_{ab}} \vec{p}_b + \vec{p}_a \cdot \vec{r}_{ab} \frac{1}{r_{ab}^3} \vec{r}_{ab} \cdot \vec{p}_b \right| \phi_{\text{el}} \right\rangle \right] \\ &+ \frac{2\pi}{3} \sum_I Z_I \frac{r_{\text{ch}}^2(I)}{\chi_e^2} \left\langle \phi_{\text{el}} \left| \sum_a \delta(\vec{r}_{aI}) \right| \phi_{\text{el}} \right\rangle \end{aligned}$$

Leading radiative correction

$$E(\alpha) = (E_0 + E_{\text{AD}} + E_{\text{NA}}) + \alpha^2 E_{\text{REL}} + \alpha^3 E_{\text{QED}} + \alpha^4 E_{\text{HQED}}$$

$$\begin{aligned} \mathcal{E}_{\text{QED}} = & \sum_{a>b} \left\{ \left[\frac{164}{15} + \frac{14}{3} \ln \alpha \right] \langle \phi_{\text{el}} | \delta(\vec{r}_{ab}) | \phi_{\text{el}} \rangle \right. \\ & \left. - \frac{14}{3} \left\langle \phi_{\text{el}} \left| \frac{1}{4\pi} P\left(\frac{1}{r_{ab}^3}\right) \right| \phi_{\text{el}} \right\rangle \right\} \\ & + \sum_a \left[\frac{19}{30} + \ln(\alpha^{-2}) - \ln k_0 \right] \sum_l \frac{4Z_l}{3} \langle \phi_{\text{el}} | \delta(\vec{r}_{al}) | \phi_{\text{el}} \rangle \end{aligned}$$

$$\ln k_0 = \frac{\langle \phi_{\text{el}} | \vec{\nabla} (H_0 - \mathcal{E}_0) \ln[2(H_0 - \mathcal{E}_0)] \vec{\nabla} | \phi_{\text{el}} \rangle}{\langle \phi_{\text{el}} | \vec{\nabla} (H_0 - \mathcal{E}_0) \vec{\nabla} | \phi_{\text{el}} \rangle}$$

Higher order QED corrections

$$E(\alpha) = (E_0 + E_{\text{AD}} + E_{\text{NA}}) + \alpha^2 E_{\text{REL}} + \alpha^3 E_{\text{QED}} + \alpha^4 E_{\text{HQED}}$$

- very difficult to calculate, known only for hydrogen and helium ⁶
- approximate formula

$$\mathcal{E}_{\text{HQED}} \approx \sum_l 4\pi Z_l^2 \left(\frac{139}{128} + \frac{5}{192} - \frac{\ln 2}{2} \right) \sum_a \langle \phi_{\text{el}} | \delta(\vec{r}_{al}) | \phi_{\text{el}} \rangle$$

⁶K. Pachucki, *Phys. Rev. A* **74**, 062510 (2006);

Variational approach in the exponential basis

Wave function expanded in terms of exponential functions

$$\phi = \sum c e^{-\alpha r_{1A} - \beta r_{1B} - \gamma r_{2A} - \delta r_{2B}} r_{1A}^i r_{1B}^j r_{2A}^k r_{2B}^l r_{12}^n$$

with

$$i + j + k + n + l \leq \Omega$$

Matrix elements with this basis functions have been investigated for many years, and are considered to be very difficult in the accurate evaluation.

Integrals

master integral

$$f(r) = \int \frac{d^3 r_1}{4\pi} \int \frac{d^3 r_2}{4\pi} \frac{e^{-u_3 r_{1A}}}{r_{1A}} \frac{e^{-u_2 r_{1B}}}{r_{1B}} \frac{e^{-w_2 r_{2A}}}{r_{2A}} \frac{e^{-w_3 r_{2B}}}{r_{2B}} \frac{r}{r_{12}}$$

others by recurrence relations

$$f_n(i, j, k, l; r) = \int \frac{d^3 r_1}{4\pi} \int \frac{d^3 r_2}{4\pi} \frac{e^{-u_3 r_{1A}}}{r_{1A}^{1-i}} \frac{e^{-u_2 r_{1B}}}{r_{1B}^{1-j}} \frac{e^{-w_2 r_{2A}}}{r_{2A}^{1-k}} \frac{e^{-w_3 r_{2B}}}{r_{2B}^{1-l}} \frac{r}{r_{12}^{1-n}}$$

→ analytic solution for arbitrary powers

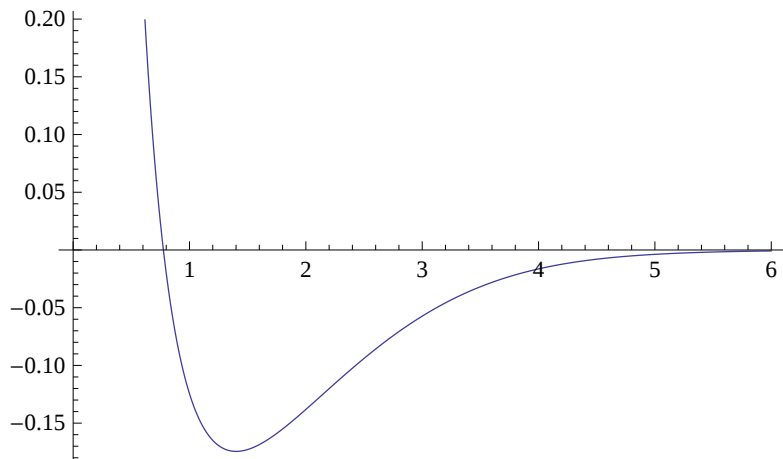
Special case: symmetric

$$f(n_1, n_2, n_3, n_4, n_5; r) = \int \frac{d^3 r_1}{4\pi} \int \frac{d^3 r_2}{4\pi} \frac{e^{-u r_{1A}}}{r_{1A}} \frac{e^{-u r_{1B}}}{r_{1B}} \frac{e^{-w r_{2A}}}{r_{2A}} \frac{e^{-w r_{2B}}}{r_{2B}} \frac{r}{r_{12}^{1-n_1}} \\ (r_{1A} - r_{1B})^{n_2} (r_{2A} - r_{2B})^{n_3} (r_{1A} + r_{1B})^{n_4} (r_{2A} + r_{2B})^{n_5}$$

$$f(0, 0, 0, 0, 0; r) = \frac{1}{4 u w} \left[e^{r(u+w)} \text{Ei}(-2 r(u+w)) + e^{-r(u+w)} \left(\gamma + \ln \frac{2 r u w}{u+w} \right) \right. \\ \left. - e^{r(u-w)} \text{Ei}(-2 r u) - e^{r(w-u)} \text{Ei}(-2 r w) \right]$$

$$f(2, 0, 0, 0, 0; r) = \frac{(u^2 + w^2)}{6 u^2 w^2} f(0, 0, 0, 0, 0; r) + \frac{e^{-r(u+w)} r^2}{12 u w} \\ + \frac{r}{24 u^2 w^2} \left[(u+w) e^{-r(u+w)} - (u-w) e^{r(u-w)} \text{Ei}(-2 r u) \right. \\ \left. - (w-u) e^{r(w-u)} \text{Ei}(-2 r w) - (u+w) e^{r(u+w)} \text{Ei}(-2 r(u+w)) \right. \\ \left. + (u+w) e^{-r(u+w)} \left(\gamma + \ln \frac{2 r u w}{u+w} \right) \right]$$

Born-Oppenheimer potential



$$E(1.4011) = -1.174\,475\,931\,400\,216\,7(3)$$

Dissociation energy of H_2 and D_2 in cm^{-1}

		H_2	D_2
α^0	Nonrelativistic	36118.7978(2)	36749.0910(2)
α^2	Mass-velocity	4.4273(2)	4.5125(2)
	1-el. Darwin	-4.9082(2)	-4.9873(2)
	2-el. Darwin	-0.5932(1)	-0.5993(1)
	Breit	0.5422(1)	0.5465(1)
	Total α^2	-0.5319(3)	-0.5276(3)
$\alpha^2 m_e/m_p$	Estimate	0.0000(4)	0.0000(2)
α^3	1-el. Lamb shift	-0.2241(1)	-0.2278(1)
	2-el. Lamb shift	0.0166(1)	0.0167(1)
	Araki-Sucher	0.0127(1)	0.0128(1)
	Total α^3	-0.1948(2)	-0.1983(2)
$\alpha^3 m_e/m_p$	Estimate	0.0000(2)	0.0000(1)
α^4	One-loop term	-0.0016(8)	-0.0016(8)
Total theory		36118.0695(10)	36748.3633(9)
Expt. 2004	Eyler	36118.062(10)	36748.343(10)
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Excitation energies of H₂ and D₂

		$J = 0 \rightarrow 1$	$v = 0 \rightarrow 1$
α^0	Born-Oppenheimer	118.55558(2)	4163.4035(1)
	Adiabatic	-0.06365(4)	-1.4029(1)
	Nonadiabatic	-0.00667(8)	-0.8365(2)
	Total α^0	118.48526(9)	4161.1641(2)
α^2	Mass-velocity	0.02713(4)	0.5347(2)
	1-el. Darwin	-0.02383(4)	-0.4994(2)
	2-el. Darwin	-0.00160(2)	-0.0391(1)
	Breit	0.00088(2)	0.0279(1)
	Total α^2	0.00258(6)	0.0235(3)
α^3	1-el. Lamb shift	-0.00109(2)	-0.0231(1)
	2-el. Lamb shift	0.00004(1)	0.0011(1)
	Araki-Sucher	0.00002(1)	0.0007(1)
	Total α^3	-0.00103(3)	-0.0213(2)
α^4	One-loop term	-0.00001(1)	-0.0002(2)
Total theory		118.48680(11)	4161.1661(5)
Experiment Jennings, Stanke		118.48684(10)	4161.1660(3)

Comparison of theoretical and experimental results for D_0 in cm^{-1} of HD

Component	D_0	δ
This work	36 405.7828(10)	
Theory		
Stanke (2009)	36 405.7814 ^a	−0.0014
Wolniewicz (1995)	36 405.787	0.004
Kołos and Rychlewski (1993)	36 405.763	−0.020
Kołos, Szalewicz, Monkhorst (1986)	36 405.784	0.001
Wolniewicz (1983)	36 405.73	0.05
Bishop and Cheung (1978)	36 405.49	−0.29
Experiment		
Merkt & Ubachs (2010)	36 405.78366(36)	0.0009
Zhang (2004)	36 405.828(16)	0.045
Balakrishnan (1993)	36 405.83(10)	0.05
Eyler and Melikechi (1993)	36 405.88(10)	0.10
Herzberg (1970)	36 406.2(4)	0.4

^a Their $D_0 = 36\,405.9794\text{ cm}^{-1}$ has been augmented by a sum of our α^3 and α^4 QED corrections equal to -0.1980 cm^{-1} .

ΔE between rotational levels of the ground vibrational state ($v = 0$) of HD, in cm^{-1}

Component	$\Delta E(0 \rightarrow 1)$	$\Delta E(0 \rightarrow 2)$
BO	89.270 629	267.196 840
Adiabatic correction	-0.036 086	-0.107 842
Nonadiabatic correction	-0.007 782(6)	-0.023 287(19)
α^0 subtotal	89.226 761(6)	267.065 711(19)
α^2 correction	0.001 948(2)	0.005 813(5)
$\alpha^0 + \alpha^2$ subtotal	89.228 709(6)	267.071 524(20)
α^3 correction	-0.000 771(1)	-0.002 303(2)
α^4 correction	-0.000 007(4)	-0.000 018(9)
Total	89.227 933(8)	267.069 205(22)
Evenson (1988)/Stoicheff (1957)	89.227 950(5)	267.086(10)

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Evenson (1988)/Stoicheff (1957)	89.227 950(5)	267.086(10)
Drouin (2010)	89.227 9326(3)	

Summary

- analytic formulae for two-center Slater integrals
- highest accuracy ever achieved for BO potential of H_2
- very good convergence with exponential basis set
- higher order QED corrections can be evaluated as accurately as for H_2^+
- more problematic are the finite nuclear mass corrections
- possible generalizations to the 3- and more electrons molecule