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Analytic approximation for electron-nucleus bremsstrahlung emissivity in the NS crust of any composition

User's guide

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We approximate the generalized Coulomb logarithm L as a function of five arguments: nuclear charge number Z , mass number of bounded nucleons A , proton core radius R_p , number density of ions (atomic nuclei) n_i , and temperature T .

Let us define suitable dimensionless parameters:

$$\Gamma = \frac{Z^2 e^2}{kT} \left(\frac{4\pi n_i}{3} \right)^{1/3} = 5.798 \frac{Z^2}{T_9} n_{i\ 37}^{1/3}, \quad (1)$$

$$\tau = \frac{kT}{\hbar Z e} \sqrt{\frac{A m_u}{4\pi n_i}} = 0.03135 \frac{T_9}{Z} \sqrt{\frac{A}{n_{i\ 37}}}, \quad (2)$$

$$\xi = R_p (3\pi^2 Z n_i)^{1/3} = \left(\frac{9\pi}{4} Z \right)^{1/3} \frac{R_p}{R_{ws}} = 0.6665 R_{p\ \text{fm}} (Z n_{i\ 37})^{1/3}. \quad (3)$$

Here $T_9 = T/(10^9 \text{ K})$, $n_{i\ 37} = n_i/(10^{37} \text{ cm}^{-3})$, $R_{p\ \text{fm}} = R_p/(1 \text{ fm})$; R_{ws} is a Wigner-Seitz cell radius.

The Coulomb logarithm is expressed as

$$L = \int_0^1 |V_{\text{eff}}(y)|^2 R_T(y) y^2 dy, \quad |V_{\text{eff}}(y)|^2 = \frac{S_{\text{eff}}(y) |F(2\xi y)|^2}{y^4}, \quad (4)$$

where $y = q/2p_F$ is a dimensionless momentum transfer (p_F being the electron Fermi momentum); $R_T(y)$ is the thermal function,

$$R_T(y) = 1 + \frac{2y^2 \ln y}{1 - y^2} \approx 1 - y, \quad \max \text{ abs error} = 0.066; \quad (5)$$

$V_{\text{eff}}(y)$ is the Fourier image of the effective potential of electron-nucleus scattering for eZ-bremsstrahlung process, the main subject of our study. Its constituents are: $F(2\xi y)$ = the nuclear form factor,

$$F(u) = 3 \frac{\sin u - u \cos u}{u^3} \approx e^{-\alpha u^2}, \quad \alpha = 0.23, \quad \max \text{ abs error} = 0.032; \quad (6)$$

$1/y^4$ = squared Fourier transform of the Coulomb potential; $S_{\text{eff}}(y)$ = the effective structure factor,

$$S_{\text{eff}} = \left(e^{w_1 y^2} - 1 \right) e^{-w y^2}, \quad (7)$$

where w_1 and w are

$$w_1 = B \frac{b\tau_{\text{eff}}}{\sqrt{(b\tau_{\text{eff}})^2 + u_{-2}^2 \exp \left\{ -7.6 \left(\tau_{\text{eff}} + \frac{18.0}{\Gamma_{\text{eff}}} \right) \right\}}}, \quad (8)$$

$$w = B \left(\frac{u_{-1}}{2u_{-2}\tau_{\text{eff}}} \exp \left\{ \tau_{\text{eff}} + \frac{216}{\Gamma_{\text{eff}}} \right\} + 1 \right), \quad (9)$$

$$B = \frac{3}{5} \frac{(12\pi^2)^{1/3}}{1 + \frac{\xi}{Z} \frac{\Gamma_{\text{eff}}}{200 + \Gamma_{\text{eff}}}} \frac{u_{-2} Z^{4/5}}{\sqrt{\Gamma_{\text{eff}}^2 + \frac{1037}{(1 + \frac{\Gamma_{\text{eff}}}{204})^4} \frac{\Gamma_{\text{eff}}^{1/2(1+0.15\xi)}}{Z^{0.10}}}}. \quad (10)$$

Here $b = 231$, $u_{-1} = 2.798$, $u_{-2} = 12.972$, and τ_{eff} and Γ_{eff} are non-linearly scaled τ and Γ , respectively. They are related to original Γ , τ , ξ and Z via the expressions:

$$\tau_{\text{eff}} = 0.095 \left(\frac{2\tau}{0.095\sqrt{\tau^2 + 4}} \right)^{\frac{1}{1 + (\frac{Z}{35.5})^2 (1 + \frac{\Gamma}{223Z})^{-1}}}, \quad (11)$$

$$\Gamma_{\text{eff}} = \Gamma \exp \left\{ - \frac{\tau_{\text{eff}}^{0.053} \left(\frac{Z}{11.3} \right)^{4/9}}{1 + \tau_{\text{eff}}^{0.50 + 0.002Z\xi} \left(\frac{\Gamma}{19.3Z^{1.7}} \right)^{\frac{1.52 + 0.90(1 - \tau_{\text{eff}})}{1 + 0.50\xi}}} \right\} \quad (12)$$

Now the Coulomb logarithm (4) is integrated analytically:

$$L = f(w + 4\alpha\xi^2) - f(w - w_1 + 4\alpha\xi^2), \quad (13)$$

$$f(x) = \sqrt{\frac{\pi}{x}} \text{erf}(\sqrt{x}) + \ln x - \text{Ei}(-x). \quad (14)$$

Here w and w_1 are given by Eqs. (9) and (8); α is defined by (6). A useful advice is to use an asymptotic value of f at $x \sim 0$:

$$f(x) \xrightarrow{x \rightarrow 0} 2 - \gamma, \quad (15)$$

where $\gamma = 0.557216\dots$ is the Euler's constant. It helps to avoid computational errors.

Finally, the neutrino emissivity takes form

$$Q = 5.362 \times 10^{18} Z^2 n_i {}_{37}T_9^6 L R_{\text{NB}} \left[\frac{\text{erg}}{\text{cm}^3 \text{ s}} \right], \quad (16)$$

where $R_{\text{NB}} = 1 + 0.00554Z + 0.0000737Z^2$ is a non-Born correction.

The details of the eZ neutrino bremsstrahlung formalism can be found in Kaminker *et al.* *Astron. Astrophys.* **343**, 1009 (1999).