Accretion heated atmospheres of X-ray bursting neutron stars

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X-ray bursting neutron stars

- X-ray bursting NSs LMXBs with thermonuclear explosions at the neutron star surface
- Sometimes close to the Eddington limit during the burst (photospheric radius expansion (PRE) bursts)
- Burst duration ~10 1000 sec

Ideal sources for NS masses and radii investigations (important for EOS!!!)



Low Mass X-ray Binary (artist veiw)



4U 1724-307 in Terzan 2

Figure from Molkov et al (2000)



Spectrum formation of X-ray bursting NSs

Compton scattering is important !!!



Model dependences

Computed using undisturbed models of hot NS atmospheres



Suleimanov et al. 2011

Suleimanov et al. 2017

Ultracompact system 4U 1820-30

Persistent spectra



Comparison with observations

Deviation from theory depends on the persistent flux



Low persistent flux

High persistent flux

Accretion heated atmospheres



 $\dot{m}_a(H_a, \overline{A}\,m_p, \upsilon_a, kT, \vartheta)$ - local mass accretion rate, g s⁻¹ cm⁻²

Accretion heated atmospheres Short history

Larkin (1960) - fast particle stopping in plasma

Zeldovich & Shakura (1969) – NS surface heated by fast particles

Alme & Wilson (1973) – the fist numerical computations

Bildsten et al. (1992) - analytical description, heavy ions destruction

Turolla et al. (1994) – "hot" solutions (T up to 10¹¹ K) Zampieri et al. (1995) – extension of AW73 computations Zane et al. (1998) – spectra of "hot" solutions Zane et al. (2000) – magnetized NS atmospheres heated by particles

Deufel et al. (2001) – application to hard quiescent spectra of LMXBs

Accretion heated atmospheres Method of computation

The approach and the code used for undisturbed hot NS model atmospheres computations was accepted (Suleimanov et al. 2012)

Additionally

- energy generation in the h
- ram pressure force

e heated layers
$$\frac{dH}{dm} \approx -\dot{m}_a \upsilon_a \frac{d\upsilon_a}{dm}$$

 $g_{ram} = \frac{dP_{ram}}{dm} \approx -\dot{m}_a \frac{d\upsilon_a}{dm}$

JU

- electron thermal conductivity $H_C \approx k_C T^{5/2} \frac{dT}{dm}, \quad \frac{dH_C}{dm} \neq 0$ (insignificant)

$$\frac{d\upsilon_a}{dm} \propto \upsilon_a^{-3} \rho^{-1} \frac{Z^2}{A}$$

Accretion heated atmospheres Adopted parameters

Neutron star: $M = 1.5 \text{ M}_{\odot}$, R = 12 km, $\ell = L/L_{Edd}$

$$\upsilon_{ff} = \left(\frac{2GM}{R}\right)^{1/2} = 1.8 \times 10^{10} \, cm/s \approx 0.6 \, c$$
$$kT_{vir} = \frac{GM \,\overline{A}m_p}{3R} \approx \overline{A} \, 61 \, MeV \, (6.7 \times 10^{11} \, K)$$

Accretion flow:

-

$$\begin{split} \ell_a &= L_a / L_{Edd} \quad \upsilon_a = \eta \ \upsilon_{ff}, \quad \eta < 1 \\ kT &= \chi \ kT_{vir}, \quad \chi < 1 \quad \vartheta \\ \overline{A} &= 1 \quad \text{-pure hydrogen} \quad \overline{A} = 4 \quad \text{-pure helium} \\ \overline{A} &\approx 1.3 \quad \text{-solar hydrogen/helium mix} \end{split}$$

Low luminosity accretion heated atmospheres with various particle velocities η

$$L = 0.001 L_{Edd}$$
 $L_a = 0.05 L_{Edd}$ $\chi = 0.1$ $\vartheta = 60^{\circ}$



Accretion heated atmospheres. Fixed parameters. Angular distribution of the emergent radiation



Accretion heated helium atmospheres. Fixed parameters. Various internal luminosities

$$L_a = 0.05 L_{Edd} \quad \eta = 0.75 \quad \chi = 0.03 \quad \vartheta = 60^{\circ}$$



Accretion heated helium atmospheres. Fixed parameters. Various heat luminosities

$$L = 0.2 L_{Edd} \quad \eta = 0.75 \quad \chi = 0.03 \quad \vartheta = 60^{\circ}$$



Accretion heated atmospheres with solar composition. Iron edge disappears in the spectrum

$$L_a = 0.05 L_{Edd} \quad \eta = 0.75 \quad \chi = 0.03 \quad \vartheta = 60^{\circ}$$



Accretion heated helium atmospheres.

Heated atmospheres less prominent evolution of the model curves at high relative velocities. The color correction values are large for heated atmospheres (corresponding dilution factors are smaller).



 $L_a = 0.05 L_{Edd} \quad \eta = 0.75 \quad \chi = 0.03 \quad \vartheta = 60^{\circ}$

Accretion heated atmospheres with solar composition. Heated atmospheres mimic atmospheres of low metal abundances



 $L_a = 0.05 L_{Edd} \quad \eta = 0.75 \quad \chi = 0.03 \quad \vartheta = 60^{\circ}$

Burster in ultracompact system 4U 1820 -- 30 Both model curves have the same fitting parameters $F_{Edd} = 0.6 \ (10^{-7} \ erg \ s^{-1}cm^{-2})$ $\Omega = R^2(1+z)^2/D^2 = 500 \ km^2/(10 \ kpc)^2$



"Clocked" burster GS 1826 -- 24

Its atmosphere has to have solar chemical composition (Heger et al. 2007). But the observed curve $K^{-1/4} - F_{BB}$ doesn't show depression at $L \approx 0.1 L_{Edd}$ typical for model curves computed for undisturbed atmospheres.



Figures from Zamfir et al. 2012

"Clocked" burster GS 1826 -- 24

But the observed curve well fitted with the "heated" model curves



Figures from Zamfir et al. 2012

Conclusions

A method of accretion heated NS atmospheres was developed.

A number of heated atmospheres was computed and the dependences on the input parameters were investigated.

The "observed" spectra of heated atmospheres of X-ray bursting NSs are well fitted with diluted blackbody.

Color correction factors are large for heated atmospheres, and dilution factors are smaller.

Iron edge disappears in the heated NS atmospheres with solar abundance of heavy elements.

Model curves $w - w f_c^4 L / L_{Edd}$ are well fitting the observed curves $K - F_{BB}$ at the later phases of the X-ray bursts.