Low-level accretion onto highly magnetized neutron stars



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X-ray pulsar

Rotating Neutron Star in binary systems





Neutron star parameters:

X-ray pulsar



$$r_{\rm co} = \left(\frac{GM}{\Omega^2}\right)^{1/3}$$

Keplerian and stellar-rotation frequencies are equal

$$r_{\rm A} = \left(\frac{\mu^4}{2GM\dot{M}^2}\right)^{1/7}$$

Alfven radius: magnetic pressure equals to the ram pressure of gas in spherical free-fall from infinity

 $r_{\rm m} = \xi r_{\rm A}$



Propeller effect



Patterson, 1994

"Propeller effect"

Illarionov & Sunyaev, 1975

- $R_m < R_c$ accretion is possible
- $R_m > R_c$ accretion is prohibited due to centrifugal barrier

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Propeller effect



 $L_{\rm lim}(R) \simeq \frac{GMM_{\rm lim}}{R} \simeq 4 \times 10^{37} \xi^{7/2} B_{12}^2 P^{-7/3} M_{1.4}^{-2/3} R_6^5 \,\rm{erg s}^{-1}$

Observational manifestation

V 0332+53

SAX J1808.4-3658

 $L_{lim} = 5 \times 10^{35} \text{ erg/s}$

GRO J1744-28



Stella et al., 1986



Patruno et al., 2016 Campana et al., 2008

Observations of the propeller effect

What is needed:

- wide range of the mass accretion rate;
- independently measured magnetic field (desired);
- sensitive and flexible X-ray telescope.

4U 0115+63, V 0332+53 and SMC X-2 in 2015



Swift/BAT + XRT monitoring

4U 0115+63, V 0332+53 and SMC X-2 in 2015



Propeller in action



Lutovinov et al., 2017

LETTER

L~100 L_{Edd} for 1.4 M_{sun} mass object!

doi:10.1038/nature13791

An ultraluminous X-ray source powered by an accreting neutron star

M. Bachetti^{1,2}, F. A. Harrison³, D. J. Walton³, B. W. Grefenstette³, D. Chakrabarty⁴, F. Fürst³, D. Barret^{1,2}, A. Beloborodov⁵, S. E. Boggs⁶, F. E. Christensen⁷, W. W. Craig⁸, A. C. Fabian⁹, C. J. Hailey¹⁰, A. Hornschemeier¹¹, V. Kaspi¹², S. R. Kulkarni³, T. Maccarone¹³, J. M. Miller¹⁴, V. Rana³, D. Stern¹⁵, S. P. Tendulkar³, J. Tomsick⁶, N. A. Webb^{1,2} & W. W. Zhang¹¹



In M82 Galaxy

P_{spin}=1.37 s



M82 as seen by Chandra



Tsygankov et al., 2016b

M82 X-2 intensity distribution



Distribution is bimodal

Tsygankov et al., 2016b

Propeller in action



GRO J1008-57 in 2016



 $P_{spin} = 94 \text{ s}, \quad E_{cyc} \sim 80 \text{ keV}, \quad d = 5.8 \text{ kpc}$

GRO J1008-57 in 2016



Thermal-viscous instability

Partial ionization of hydrogen at ~6500 K cause an abrupt change in opacity and viscosity making the disc locally unstable. Critical accretion rate, above which the disc is stable:



Thermal-viscous instability

Partial ionization of hydrogen at ~6500 K cause an abrupt change in opacity and viscosity making the disc locally unstable. Critical accretion rate, above which the disc is stable:

$$\dot{M} > \dot{M}_{\rm hot} \approx 6 \times 10^{16} r_{\rm out,10}^3 \ {\rm g \, s}^{-1}$$

Or below:

$$\dot{M} < \dot{M}_{cold} \simeq 3.5 \times 10^{15} r_{10}^{2.65} M_{1.4}^{-0.88} \text{ g s}^{-1}$$

Lasota, 1997

Substituting the magnetospheric radius instead of *r*:

$$L < L_{\text{cold}} = 9 \times 10^{33} k^{1.5} M_{1.4}^{0.28} R_6^{1.57} B_{12}^{0.86} \text{ erg s}^{-1}$$

GRO J1008-57 in 2016



Propeller effect vs cold disc

The final state of the source after an outburst is determined by two fundamental parameters of the neutron star: magnetic field and spin period. Equating the expressions for luminosities L_{cold} and L_{prop} one can derive the critical value of the spin period as a function of the neutron star magnetic field:

$$L_{\rm lim}(R) \simeq \frac{GM\dot{M}_{\rm lim}}{R} \simeq 4 \times 10^{37} \xi^{7/2} B_{12}^2 P^{-7/3} M_{1.4}^{-2/3} R_6^5 \,\rm erg \,\, s^{-1}$$

 $L_{\rm cold} = 9 \times 10^{33} k^{1.5} M_{1.4}^{0.28} R_6^{1.57} B_{12}^{0.86} \, {\rm erg \, s}^{-1}$

$$P^* = 36.6 \, k^{6/7} \, B_{12}^{0.49} \, M_{1.4}^{-0.17} \, R_6^{1.22} \quad \text{s}$$

Tsygankov et al., 2017

Propeller effect vs cold disc



Conclusion (I)



Conclusion (II)

 $P^* = 36.6 k^{6/7} B_{12}^{0.49} M_{14}^{-0.17} R_6^{1.22}$ S



Stable accretion from the cold disc vs propeller effect